# PROJECT SQUID

**TECHNICAL REPORT MIT-50-PU** 

VISCOSITY OF NITROGEN, HELIUM, HYDROGEN, AND ARGON FROM -100° TO 25°C UP TO 150-250 ATMOSPHERES

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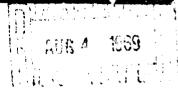
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JET PROPULSION CENTER
SCHOOL OF MECHANICAL ENGINEERING
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LAFAYETTE, INDIANA

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## PROJECT SQUID

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Viscosity of Nitrogen, Helium, Hydrogen, and Argon from -100° to 25°C up to 150-250 Atmospheres\*

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The range of a previously developed capillary-flow viscometer has been extended to -100°C, and the techniques of obtaining and determining the attainment of steady state have been improved. The viscosities of nitrogen, helium, hydrogen, and argon have been measured by an absolute method from -100° or -90° to 25°C and up to 150-175 atm (250 atm for N<sub>2</sub> at 25°C). The accuracy is estimated to be 0.1 - 0.2%, but is somewhat worse for Ar at -100°C. A number of empirical and theoretical analyses of the viscosity-density dependence have been made; while inconclusive, the evidence tends not to support the presence of a logarithmic term in the density expansion. Correlations of the second and the third virial coefficients of viscosity have been made.

We report here a series of absolute measurements of the viscosity of four gases as a function of temperature and density. The purpose of this work is to obtain experimental data on this transport property in regions where such information is unavailable or insufficient; to investigate the temperature dependence of the virial coefficients for the viscosity; and to provide, if possible, a test of proposed theories on the density expansion.

We have in previous publications<sup>1,2</sup> described our viscometer and given results for a number of gases at temperatures from -50° to  $150^{\circ}$ C. The present work extends the temperature range to -90°C for N<sub>2</sub> and He and to -100°C for H<sub>2</sub> and Ar, and repeats the previous measurements at -50° and  $25^{\circ}$ C for more closely spaced values of the density. As before, the upper pressure limits are in the range 150-175 atm, except for N<sub>2</sub> at  $25^{\circ}$ C, which we studied to about 250 atm. The precision of the data is better than 0.1%, and the results should be accurate to 0.1 - 0.2%, with the exception of Ar at -100°C (see discussion below).

At the time of publication of Ref. 1, there were hardly any other published measurements of the viscosity of gases below 0°C and above 1 atm, and none of comparable accuracy. This situation has improved in recent years, especially with Kao and Kobayashi's work on N<sub>2</sub> and He from -90° to 50°C and Diller's on p-H<sub>2</sub> below 100°K, as well as some extensive measurements of much lesser accuracy mentioned below. However, none of these papers contains enough low-density data to give good values of the second virial coefficient of viscosity which is of theoretical interest. 5-9 For this reason the present measurements have been extended to lower pressures than our previous work. 1,2

#### APPARATUS AND PROCEDURES

The apparatus and technique of measurement have been described in Ref. 1, and only some changes need be noted here; for a more extensive discussion, see Ref. 10. The apparatus is a constant-volume capillary-flow viscometer, in which an injector at 25°C drives gas at a known flow rate through a capillary at the experimental temperature, the resulting pressure difference being determined with a differential manameter.

The only major change is the addition of a fused-quartz precision pressure gage (Texas Instruments Inc., Houston, Texas) in parallel with the differential manometer. This instrument registers a pressure difference by deflection of a quartz Bourdon tube. We had hoped to use it for absolute measurements, but found that the variation of the index of refraction of a gas with pressure affects the optical path in the instrument, so that its calibration varies with pressure. We therefore continued to use the mercury manometer, but made use of the pressure gage to monitor the approach to steady state.

Steady state is defined here as invariance with time of the pressure difference across the capillary during a run. Strictly speaking, such steady-state behavior is impossible in a constant-volume apparatus, since the pressure difference steadily increases with time; 11 however, we have shown 10 that this increase is negligible for the dimensions of our apparatus. A more serious problem is the relaxation of the transient process set in motion when a run begins: in particular, the rapid imposition of the pressure difference is an essentially adiabatic process, producing a temperature difference across the capillary. For most of the measurements reported here, we in fact observed the attainment of steady state (constant within c. 0.001

in. Hg) with the pressure gage as noted above. In addition, a rough but adequate model of the transient behavior has been derived  $^{10}$  and approximately confirmed experimentally, to the extent that we believe steady state to have been reached in virtually all the previous measurements with this apparatus. (This applies even for injector flow rates up to 0.05 cm $^3$ /sec, although we now work only at  $\leq 0.03$  cm $^3$ /sec.)

Three capillaries were used in the present work, all about 36 cm long and 0.019 cm in diameter. The last two were calibrated by the electrical-resistance method, which we have found to be both more convenient and more precise than the gravimetric method. More careful analysis of our previous techniques with the latter method led us to the conclusion that the viscosity values in Refs. 1 and 2 should be decreased by about 0.12.

For each capillary the value of the kinetic-energy constant m was determined in the same way as before , except that individual runs were weighted in accordance with their a priori precision. We again took an upper Reynolds-number limit of 1600, below which we found no correlation of m with Reynolds number or other variables. For the three capillaries the values of m obtained were  $1.199 \pm 0.013$ ,  $1.162 \pm 0.016$ , and  $1.199 \pm 0.025$ ; these values were used to calculate the viscosities in Table I.

## EXPERIMENTAL RESULTS

Our measured viscosities are given in Table I and plotted in Figs. 1-4; the deviations listed in the table indicate primarily the scatter of the measurements at each pressure. The paragraphs which follow give our density sources and some discussion of the results. More extensive comparisons with previous work can be found in Ref. 10.

## Nitrogen

Nitrogen<sup>14</sup> was studied up to 253 atm at 25°C; at -50° and -90°C, as for the other gases, limits were chosen in the range 150-175 atm. The results are shown in Fig. 1. The densities used were those of Michels et al. <sup>15</sup> at 25°C and of Canfield et al. <sup>16</sup> at lower temperatures. Our results are in good agreement <sup>17</sup> with our previous work<sup>1</sup>, and thus are still about 0.4% lower than those of Kestin and Leidenfrost <sup>18</sup> and Michels and Gibson <sup>19</sup> at 25°C and moderate pressures. Agreement with Kao and Kobayashi at -50° and -90°C is within their precision. On the other hand, our disagreement with Ross and Brown <sup>20</sup> is as reported before, <sup>1</sup> and Filippova and Ishkin's <sup>21</sup> measurements deviate from ours in the same way; Goldman's <sup>22</sup> data at -61° and -78.5°C are 1-6% higher than values interpolated from ours; and extrapolation of our -90°C curve indicates that Zhdanova's <sup>23</sup> results for  $\rho \geqslant 0.38$   $\because$ cm<sup>3</sup> are as much as 20% higher.

#### Helium

Helium was also studied at 25°, -50°, and -90°C; the results are shown in Fig. 7. The densities are based on the work of Miller et al. 25°C and of Canfield et al. 10 elsewhere. We are again in good agreement with

our previous work<sup>1</sup> and about 0.3% lower than the results of Kestin et al.<sup>18</sup> around 25°C. At -50° and -90°C, Kao and Kobayashi's<sup>3</sup> results average 0.2-0.3% lower than ours. Golubev and Gnezdilov<sup>26</sup> obtain viscosities consistently steeper in the density than all the above work, but closely resembling those of Ross and Brown<sup>20</sup>, which we regard as unreliable.<sup>1</sup> Kestin's results, Kao and Kobayashi's, and our's all tend to confirm the existence of very shallow minima in the viscosity-density curves.

## Hydrogen

Hydrogen<sup>27</sup> was studied at 25°, -50°, and -100°C; the results are shown in Fig. 3. The densities were all taken from Michels et al. <sup>28</sup> Agreement with our previous work<sup>2</sup> was excellent at 25°C and fairly good (the present data averaging slightly higher) at -50°C. Near 25°C, Kestin and Leidenfrost's results are about 1% higher than ours, but the recent work of Kestin and Yata<sup>29</sup> is only about 0.5% higher, being thus in good agreement with Michels, Schipper, and Rintoul.<sup>30</sup> There is no other high-precision work in our lower-temperature range; both Golubev and Shepeleva<sup>31</sup> and Rudenko and Slyusar <sup>32</sup> obtain results which increase with density somewhat more rapidly than ours or Michels'.

## Argon

Argon<sup>33</sup> was also studied at 2,°, -50°, and -100°C; the results are shown in Fig. 4. The densities were obtained from Michels, Wijker and Wijker<sup>34</sup> at 25°C, and from Michels, Levelt, and de Graaff<sup>33</sup> at the lower temperatures.

Agreement with our earlier work<sup>1</sup> is again fairly good. As before, at 25°C and moderate densities we are about 0.3% below Kestin et al.<sup>18</sup> and 0.2% above Michels. Botzen, and Schuurman. Filippova and Ishkin<sup>1</sup>s<sup>21</sup> low-temperature data deviate from ours in the once way as for 11. Our -100°C isotherm requires

special discussion, which must be postponed until the questions of density dependence have been considered.

### DENSITY DEPENDENCE OF VISCOSITIES

For each isotherm the viscosity was fitted to a least-squares power series in the density:

$$\eta = a + b\rho + c\rho^2 + \dots$$
 (1)

These fits were tried for polynomials of different degrees, the "best" fit (to all the points of the isotherm) being chosen by a variety of statistical tests; <sup>37</sup> the constants of the resulting series are listed in Table II. The standard deviations ( $\sigma$ ) of the fits are in each case well under 0.2%, except for N<sub>2</sub> at -90°C and Ar at -100°C (for both of which the individual viscosities also have quite high deviations), thus confirming our estimate of precision (see below).

The "best" fits just described, however, do not necessarily yield the best values of the virial coefficients, as has been shown by Hanley, McCarty, and Senjers. These authors found, by fitting series to data subsets of increasing density range, that (for example) a good linear fit could be obtained up to some limiting density  $\rho_{\ell}$ , beyond which the data deviated from linearity, but that good coefficients for a quadratic fit were not obtained below some density appreciably greater than  $\rho_{\ell}$ . Most of our isotherm ranges fall in such intermediate regions. We have therefore carried out such subset fits for our data, and the resulting estimates for the best of values of a, b, c are given in Table III for N and Ar: the

scatter was two great for satisfactory results by this technique for He and H<sub>2</sub>. The values in Table III should thus be taken as our best estimates of these virial coefficients, while those in Table II are best for interpolation over our entire density range.

For  $N_{\varphi}$  and Ar fits were also made to the series

$$\eta = a + b\rho + \rho^2 (c' + d'ln\rho), \qquad (2)$$

which has been proposed on theoretical arounds.<sup>38</sup> In each case<sup>10</sup> the goodness of fit was no better ( (usually) slightly worse than for the power series (1) with an equal number of coefficients. Our work thus offers no support for the logarithmic-term hypothesis, whereas Hanley, McCarty, and Sengers<sup>9</sup> did find some support in an analysis based primarily on thermal conductivity data.

In addition to our individual-isotherm fits, further tests were made of the hypothesis that the residual viscosity is a function of density only, i.e., that

$$\eta(T,\rho) = \eta(T,0) + \Delta\eta(\rho) . \tag{3}$$

This is an empirical correlation, the breakdown of which at low temperatures and high densities has been fairly well established. 4,23 However, where it is valid it not only is quite useful for purposes of estimation, but holds to a remarkable degree of accuracy. Such fits were therefore made, as described in Ref. 1, for each of the gases studied; the standard deviations of the "best" fits are 0.26µP for N<sub>2</sub>, 0.14 for He, 0.14 for H<sub>2</sub>, 1.41 for Ar (0.46 with -100°C omitted). (See Ref. 10 for the constants in these series and detailed comparison graphs showing to what extent the isotherms are in fact parallel.) Except for Ar, these results are quite consistent with the validity of Eq. (3) for our gases, within our experimental accuracy.

The Ar results, however, are problematical. Even the 25° and -50°C results are definitely not parallel within their own precision (the 25°C data being parallel to Kestin's and Michels' results), but the discrepancy here is only of the order of  $1 \mu P$ . The  $-100^{\circ} C$  data, in contrast, dip as much as 4% below the values predicted from Eq. (3) and Michels!  $\Delta\eta(\rho)$ , then rise more steeply at higher densities. (It may be significant that the greatest discrepancy, nearly 10 uP, is at 0.5 g/cm<sup>3</sup>, quite close to the critical density of 0.531 g/cm<sup>3</sup>.) The -50°C deviations, while much smaller, are in the same direction (as are the previous rough results at -78.5°C). As noted above, Eq. (3) is known to fail in the low-temperature liquiddensity range, but the previous evidence 4,23 indicates that the departure is in the direction of higher values of  $\Delta\eta(\rho)$ , the opposite of what we observe here; no effect of either kind appears in the present No data, at similar reduced temperatures and densities. The poor precision of the -100°C data is far too small an effect to account for this discrepancy, and further low-temperature Ar measurements are clearly desirable.

To return to the individual-isotherm fits, we have reduced the values of the second virial coefficient of viscosity, b, from Table III for  $N_2$  and Ar, Table II for  $H_2$  and He for comparison with the theories of Kim and Ross and Hoffman and Curtiss. The reduced values,  $b_{\eta}^{*}$ , are shown in Fig. 5, along with the two theoretical curves; see Refs. 8 for comparison with several other theoretical models, 5,6,39 none of which fit the experimental data well. Also shown in Fig. 5 are the values of  $b_{\eta}^{*}$  from the data of Kestin and Leidenfrost,  $b_{\eta}^{*}$  to which the parameters in Kim and Ross's theory were fitted; the present values are in general higher, but the shape of the curve remains unchanged. The theory of Kim and Ross seems reasonably

correct, except for its failure to predict the negative values of b for He; on the other hand, that of Hoffman and Curtiss agrees well with the He data but fails at lower reduced temperatures. These results can be explained as due to Kim and Ross taking inadequate account of repulsive interactions and Hoffman and Curtiss's neglect of orbiting, effects important at high and low reduced temperatures respectively. The plateau predicted by Kim and Ross for  $(1/T^*) \ge 0.3$  presumably is related to the nearly temperature independent residual viscosity discussed above, although the latter appears to apply even at densities where higher virial coefficients are dominant.

In connection with the point just made, a correlation of the third virial coefficient should be useful; to our knowledge, this has not been attempted before. While no theory for the value of c exists, the appropriate reduction is easily made by dimensional analysis: since we have 8

$$b_{\eta}^{*} \left(T^{*}\right) = \frac{\left(6/\sigma\right)}{\left(\epsilon/m\right)^{\frac{1}{2}}} b_{\eta} \qquad , \tag{4}$$

where  $\sigma$  and  $\varepsilon$  are Lennard-Jones parameters and  $T^*\equiv kT/\varepsilon$  , it is clear that a dimensionless  $c_n^*$  should be given by

$$c_{\eta}^{*}(T^{*}) = \frac{(1/\sigma^{4})}{(\epsilon/m^{3})^{\frac{1}{2}}}c_{\eta}$$
 (5)

If  $\sigma$  is given in A and  $(\varepsilon/k)$  in  ${}^{\circ}K$ , the reduction factor  $\sigma^4(\varepsilon/m^3)^{\frac{1}{2}}$  equals 54.92  $\sigma^4$   $(\varepsilon/k)^{\frac{1}{2}}$   $\mu P-cm^6/g^2$ , where M is the molecular weight. Values of  $c_{\eta}^{\phantom{\eta}}$  were therefore obtained from our data as just described for  $b_{\eta}^{\phantom{\eta}}$ , the results being plotted in Fig. 6. Once again a plateau is obtained over most of the temperature range.

## ACCURACY

Our previous estimate of experimental precision was 0.001 in.

Hg in the pressure difference. Although in our low-temperature work we were frequently restricted (by steady-state considerations) to pressure differences much less than an inch, the reproducibility of the viscosities (Table I) remained better than 0.1% in nearly all cases (even for H<sub>2</sub> at -100°C, where none of the pressure differences exceeded 0.4 in. Hg). The exceptions to the above statement are of course N<sub>2</sub> at -90°C and Ar at -100°C. As mentioned in the previous section, the standard deviations of our isotherm least-squares fits are consistent with these estimates of precision. Our average precision is thus about 0.05% near room temperature, but becomes worse than 0.1% at low temperatures and high densities (perhaps due to lack of thermal equilibrium in the apparatus).

Our accuracy should be somewhat worse than this. For our low-temperature  $N_2$  and He measurements, there is the possibility of errors larger than 0.1% in the density ratios and thus the viscosities, due to our having combined densities from two different sources; the single-source  $H_2$  and Ar density ratios should be good to about 0.02%. Combining this with other known sources of error in the calibrations and measurements, we estimate our accuracy as 0.1 - 0.2%.

However, accuracy can be established objectively only by comparison with other work, and the situation in this respect is unsettled. Our agreement with previous work with this apparatus<sup>1,2</sup> is within the limits just stated, but at -50°C there are some apparently systematic discrepancies<sup>10</sup> at the outer range of these limits. The most precise other work available

for comparison is that of Kestin et al., 18,29 done with oscillating-disc viscometers, and the discussions above have indicated that our results are consistently several tenths of a percent lower than Kestin's; the possibility of flaws in the theory of either or both instruments may be worth investigating. We are in good agreement with Kao and Kobayashi<sup>3</sup> for N<sub>2</sub> and slightly higher for He, while our discrepancies with Michels et al. 19,30,36 are not consistent (both of these being capillary measurements). There are no other dense-gas data sufficiently accurate for useful comparison, but it should be noted that our 1-atm viscosities (based on Table II), except for He, are consistently an average of about 0.5% lower than the bulk of previous work is since we deliberately extended our work to lower densities than before, this can no longer be attributed to extrapolation error. The weight of the evidence suggests the possibility of systematic errors, but is not sufficiently consistent to be conclusive.

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## FIGURE CAPTIONS

- Fig. 1 Viscosity of nitrogen versus density: present results; Flynn,

  Hanks, Lemaire, and Ross (points nearly coinciding with present

  data indicated by ); × Kao and Kobayashi.
- Fig. 2 Viscosity of helium versus density: present results; O Flynn,
  Hanks, Lemaire, and Ross; × Kao and Kobayashi.
- Fig. 3 Viscosity of hydrogen versus density: present results; O Barua,
  Arzal, Flynn, and Ross<sup>2</sup>; △ Diller. 4
- Fig. 4 Viscosity of argon versus density: present results; Flynn,

  Hanks, Lemaire, and Ross (points nearly coinciding with present

  data indicated by •).
- Fig. 5 Comparison of theoretical  $^{5,6}$  and experimental values of  $b_{\eta}^{\phantom{\eta}}$ , the reduced second virial coefficient of viscosity: present data,

    $N_2$ ,  $H_e$ ,  $H_2$ , Ar; Kestin and Leidenfrost,  $^{18}$   $N_2$ ,

    $H_2$ , Ar,  $H_2$ ,  $H_2$ , •
- Fig. 6 Values of  $c_n^*$ , the reduced third virial coefficient of viscosity, from our data:  $\bullet$   $N_2$ ,  $\blacksquare$  He,  $\blacktriangledown$   $H_2$ ,  $\blacktriangle$  Ar (from Table II for  $N_2$  and Ar, Table III for He and  $H_2$ ).

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Pressure (atm)	Density (#/cm <sup>3</sup> )	Viscosity (L.P)	Pressure (atm)	Density $(g/cm^3)$	Viscosity (µP)
	N <sub>2</sub> , 25°C			<sup>N</sup> ≥, -90°c	
6.77	٥.يعبر	177.93+0.04	5.28	0.01000	119.32+0.11
98.03	0.02396	130.2770.03	10.39	0.02000	119.65±0.06
=	•	180.14+0.02	20.16	0.04000	123.02+0.07
=	2	130.05-0.07	29.37	0.06000	127.60+0.34
ડે. છે.	26240.0	163.82-0.04	42.30	00060.0	133.46+0.13
54.20	0.06000	186.0540.07	54.38	0.12000	140.76+0.18
<b>R</b> .3	0.07188	188.7840.05	65.85	0.15000	149.55+0.25
83.71	0.09585	193.97+0.06	76.95	0.18000	159.1940.08
105.24	0.11961	200.19+0.07	87.94	0.21000	169.73.40.10
•	0.11981	200.07±0.06	99.10	0.24000	182,5040,20
123.81	0.14000	205.59+0.10	114.80	0.28000	201.06+0.22
127.33	0.14377	206.7940.06	132.26	0.32000	222.4240.13
142.71	0.16660	211.7240.11	152.63	0.36000	247.19+0.17
152.23	0.18000	218.4140.12		030 AT	l
173.88	0.19169	222.62+0.08	,	ne, 22 c	
198.96	0.21565	231.7240.09	6.13	0.001000	198.07+0.03
243.52	0.26358	251.6640.08	12.30	0.002000	197.91+0.03
	000		24.75	0.001000	197.83+0.02
	2 0C= 12m		37.35	0.00900.0	197.7640.07
6.48	0.01000	141.75+0.03	56.53	0.009000	197.84+0.03
12.85	0.02000	142.9170.11	76.07	0.012000	197.8840.03
8.3	0.04000	145.9640.14	95.98	0.015000	197.92+0.03
37.40	0.06000	149.52-0.08	116.25	0.018000	197.95+0.03
12.67	0,08000	153.68-0.11	150.90	0.021000	190.12+0.03
8.3	0.10000	158.2340.12	L7(-93	0.024000	190.2/±0.0/
72.43	0.12000	163.31+0.13			
83.92	0.14000	168.8940.08			
95.46	0.16000	175.13+0.21			
3	: (	175.03+0.17			
107.13	0.18000	181.77+0.15			
119.03	0.2000	186.9640.19			
131.26	0.22000	196.8940.21			
143.93	0.24000	204.92+0.15			
157.16	0.26000	214.11+0.13			
01.17	0.2000	663.0640.13			

Pressure (atm)	Density (g/cm <sup>3</sup> )	Viscosity (µP)	Pressure (atm)	Density (g/cm <sup>3</sup> )	Viscosity (µP)
Alternative and the state of th	He, - 50°C			H <sub>2</sub> , 25°c	
4.59	0.001000	+	50.01	0.004000	
9.21	0.002000	+	63.00	0.005000	+
18.53	0.004000	8	76.19	0.006000	0  +
27.91	0.006000	+	89.59	0.007000	0  +
42.35	000600.0	+	103.22	0.008000	0 + 170.
57.00	0.012000	+	117.08	0.009000	.50 + 0
71.93	0.015000	162.74 7 0.05	138.31	0.010500	$60.06 \pm 0.05$
87 . 14	0.018000	<del>+</del>	160.10	0.012000	0 + 14.
102.64	0.021000	8 1+1		C	
118.43	0.024000	.16 +		H <sub>2</sub> , - 50°C	
1,4.53	0.027000	+ 52.		ı	
150.93	0.030000	, 6 +	4.56	0.000500	+
167.65	0.033000		9.14	0.001000	+
	•	i	18.40	0.002000	+
	He, -90°C		27.77	0.003000	+
			37.27	0.004000	73.63 + 0.05
7.56	0.002000	+	76.90	0.005000	+
15.21	0.004000		56.67	0.006000	+
22.95	0.006000	<b>†</b>	66.58	0.007000	+
\$4.75	0.009000	+	あ.92	0.008000	1+
146.77	0.012000	142.41 7 0.07	86.85	0.009000	74.60 + 0.03
63.16	0.016000	+	102.48	0.010500	+
79.97	0.020000	1+	118.50	0.012000	+
97.21	0.024000	+	134.92	0.013500	+
114.89	0.028000	+	151.78		•
133.04	0.032000		169.10	0.016500	77.52 + 0.06
151.67	0.036000	% 1+			ı
170.78	0.040000	145.29 ± 0.05			
	H., 25°C				
<b>6.</b> 8	0.000500	+   +			
26.63	000000	88.81 + 0.03			
77.23	0.003000	1+			
		l			

Pressure (atm)	Density (r/cm <sup>3</sup> )	Viscosity (µP)	Pressure (atm)	Density $(\varepsilon/cm^3)$	Viscosity (µP)
	H2, -100°c			Ar, -50°C	
7.08	0.001000	20.0 + 72.00	45.4	0.01000	176.41 + 0.02
14.83 21.65	0.00.00	+   +	9.00 17.68	0.04000	+1+
( c)	0.024500	1+	26.06	00090.0	1+
£3.60	0.006000	1+	34.15	0.080.0	+
16.40	0.007500	1+1	41.98	0.10000	1+1
=	Ľ	+	15.64	0.12000	+
16.67	0.005000	+	ó <del>4</del> .10	0.16000	+
=	Ξ	<b> </b> +	7.7.90	0.2000	1+
(4.8)	0.610500	+	91.14	0.24000	+
90.55	0.012000	+	103.98	0.28000	1+
197.10	0.014000	1+	116.60	0.32000	1+
124.63	0.016000	<b> </b> +	129.19	0.36000	+
141.97	0.018000	1+	141.92	00007.0	1+
160.36	0.020000	ı <b>+</b>	155.02	0.44000	<b>i</b> +
	Ar, 25°C	ı	168.72	0.48000	1+1
7.30	0.01198	226.61 + 0.03			
法:法	0.02396	+			
28.85	0.04792	+			
42.95	0.07189	+			
56.87	0.09585	+			
30.05	0.11981	+1			
Q£ .35	0.14577	+			
111.60	0.19170	260.72 ± 0.10			
1,38.93	0.23962	273.40 + 0.11			
166.70	0.28754	267.24 + 0.10			

TABLE I. (continued)

Pressure (atm)	Density $(g/cm^3)$	Viscosity (µP)
	Ar, - 100°C	
6.89	0.02000	0+
13.33	0.04000	0  +
19.36	0.06000	∻
19,36	0.06000	ت 4ا
27.64	0.090.0	0  +
35.09	0.12000	0  +
43.85	00.16000	0  +
51.14	0.20000	)  +
59.56	0.25000	181.76 7 0.17
66.45	0.3000	ە + 1
72.43	0.35000	0  +
77.78	0.40000	ن ۱+
82.79	0.45000	ن ۱+
87.76	0.50000	+
93.01	0.55000	o !+
98.96	0.6000	329.45 7 0.45
106.12	0.65000	360.02 7 0.23
115.17	0.7000	396.55 7 0.15
126.96	0.75000	436.04 7 0.40
142.61	0.80000	479.25 7 0.52
158.86	0.84000	520.37 ± 0.25

 $a + bp + cp^2 + ...$ TABLE II. Constants in viscosity-density power series fits (A

3 <b>8</b> 5	(°c)	а (µР)	ь (пР-ст <sup>3</sup> /В)	c (μΡ-cm <sup>6</sup> / <sub>3</sub> <sup>2</sup> )	d μΡ-cm <sup>9</sup> /κ <sup>3</sup> )	$(\mu P-cm^{12}/\kappa^4)$	σ (πb)
7	25 -50 -90	176.96 ± 0.07 140.59 ± 0.08 117.28 ± 0.34	116.5 ± 1.4 108.4 ± 1.3 137.1 ± 9.1	634 + 5 670 + 484 670 ÷ 484	520+108		0.12 0.11 0.38
	25 -50 -90	198.01 ± 0.05 162.82 ± 0.07 142.71 ± 0.09	-41.1 ± 9.8 -35.2 ± 10.0 -46.2 ± 11.0	2175 ± 393 1932 ± 294 2835 ± 261			0.06
•	25 -50 -100	88.55 ± 0.03 72.93 ± 0.06 60.99 ± 0.08	144.8 ± 12.3 98.2 ± 18.1 223.7 ± 18.0	7693 ± 985 10531 ± 1058 6148 ± 843			0.09
	25 -50 -100	225.30 ± 0.12 175.27 ± 0.08 138.04 ± 0.58	109.0 ± 3.8 101.5 ± 0.9 86.6 ± 10.7	442 ± 31 367 ± 2 357 ± 52	-249+ 70	212 + 54	0.11 0.14 0.61

TABLE III. Best estimates for first three virial coefficients of viscosity, from fits to subsets of the present data.

Gas	T (°C)	<b>ε</b> . (μΡ)	b (µP-cm <sup>3</sup> /g)	с (µP-cm <sup>6</sup> /g <sup>2</sup> )
$N_2$	25	176.96	116.5	633
	-50	140.49	111.6	655
	-90	117.52	122.7	603
Ar	25	225.17	116.2	359
	<del>-</del> 50	175.19	103.2	363
	-100	137.83	91.0	338

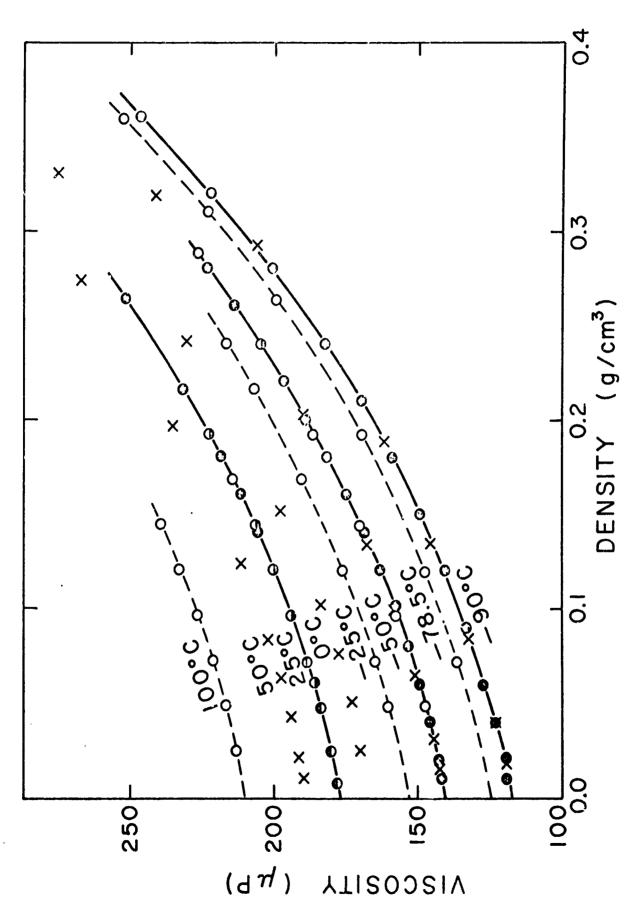
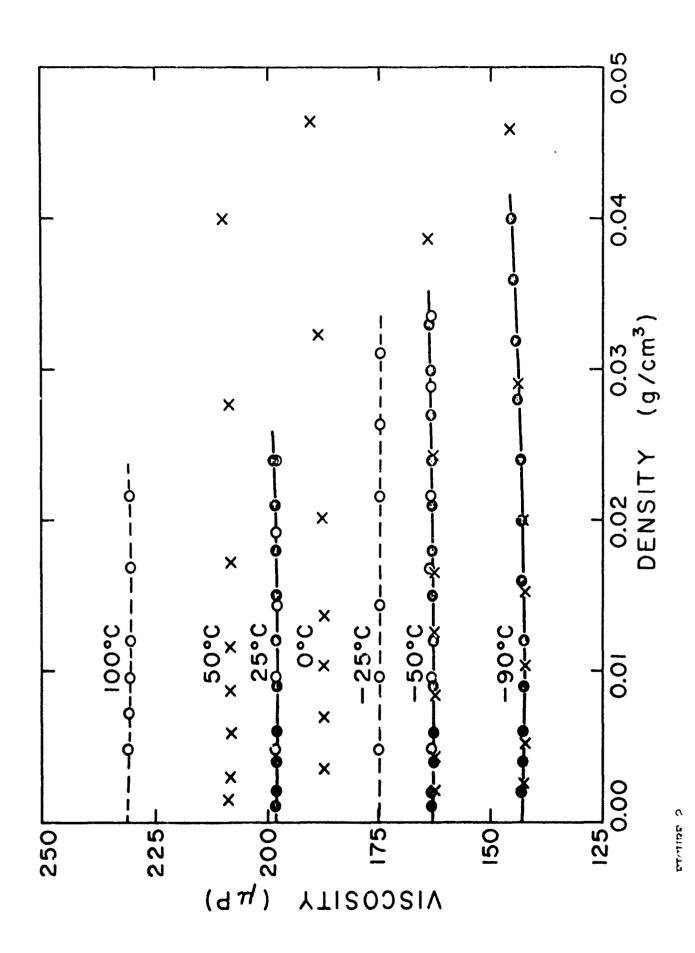
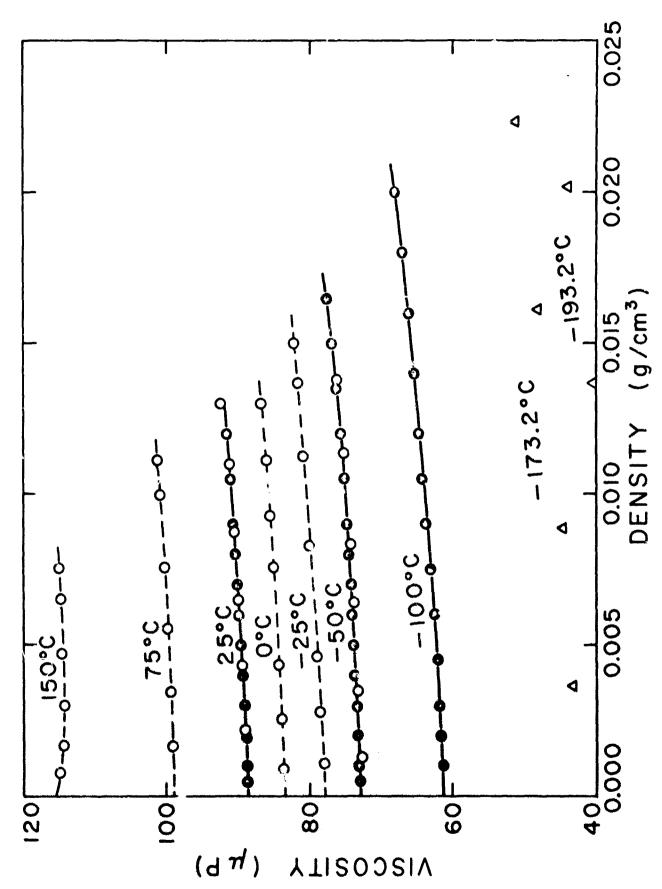
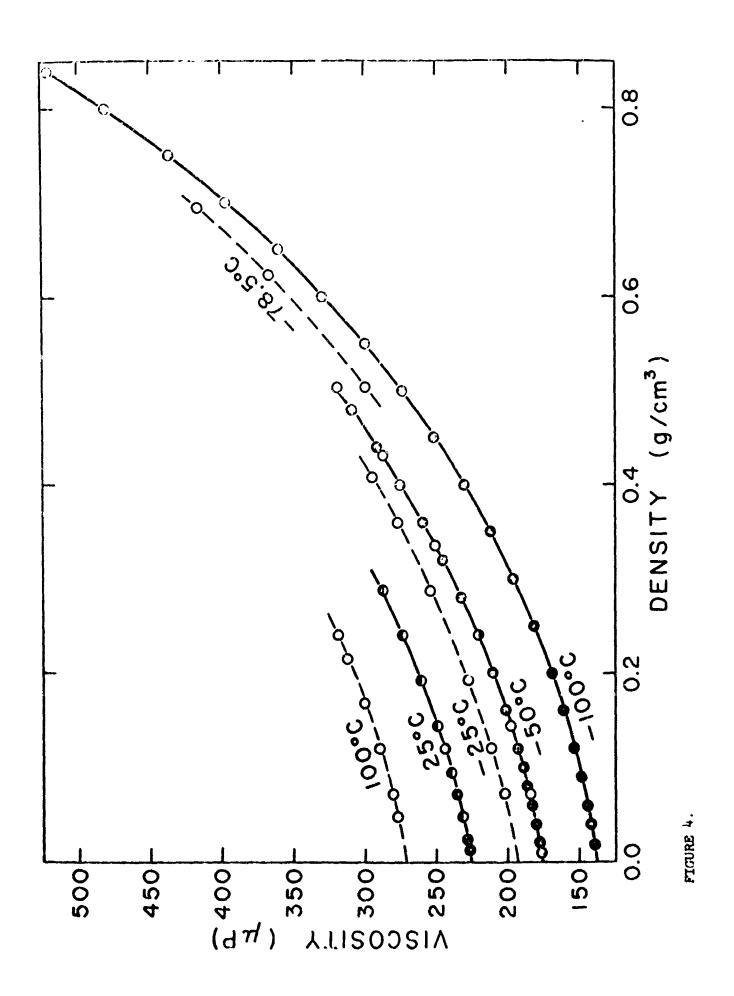


FIGURE 1.





PICURE 3.



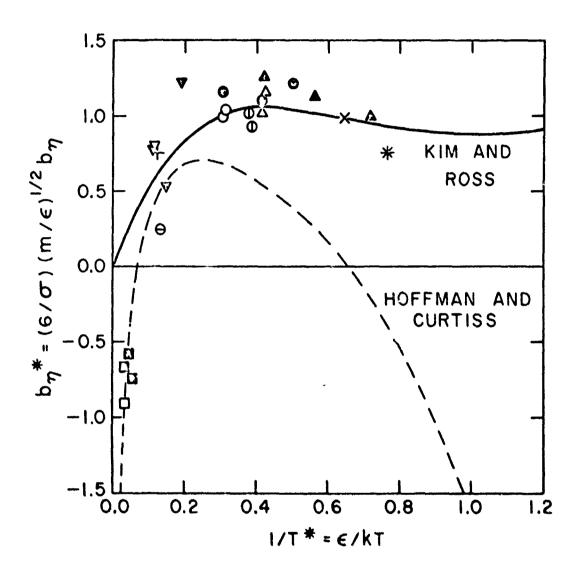


FIGURE 5.

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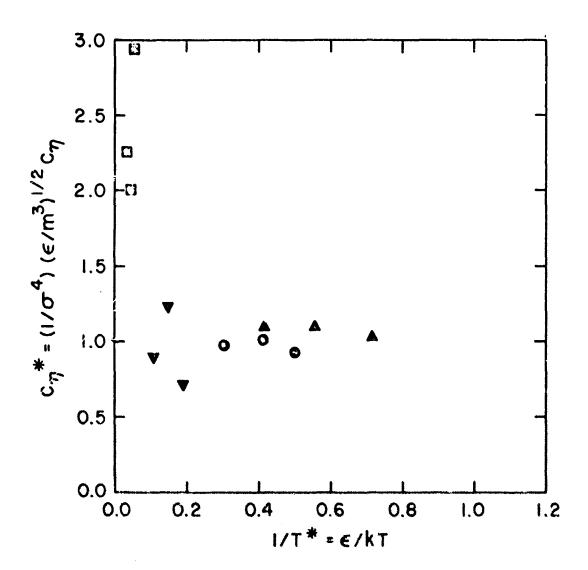


FIGURE 6.

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VISCOSITY OF NITROGEN, HELIUM, HYDROGEN, A 150-250 ATMOSPHERES	ND ARGON FRO	<b>ом -100°</b> то	25°C ÚP TO			
4. CESCRIPTIVE NOTES (Type of report and inclusive dates)  N/A			1			
J. A. Gracki, G. P. Flynn and John Ross						
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The range of a previously developed of to -100°C, and the techniques of obtaining						

The range of a previously developed capillary-flow viscometer has been extended to  $-100^{\circ}$ C, and the techniques of obtaining and determining the attainment of steady state have been improved. The viscosities of nitrogen, helium, hydrogen, and argon have been measured by an absolute method from  $-100^{\circ}$  or  $-90^{\circ}$  to  $25^{\circ}$ C and up to 150-175 atm (250 atm for N<sub>2</sub> at 25°C). The accuracy is estimated to be 0.1-0.2%, but is somewhat worse for Ar at  $-100^{\circ}$ C. A number of empirical and theoretical analyses of the viscosity-density dependence have been made; while inconclusive, the evidence tends not to support the presence of a logarithmic term in the density expansion. Correlations of the second and the third virial coefficients of viscosity have been made.

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Security Classification ROLE N<sub>2</sub>, He, H<sub>2</sub>, A viscosity-density dependence virial coefficients low temperature high pressure viscometer capillary flow Best Available Copy

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